

Uniform scaling of electronic coords (DFT-1)

$$\bar{r}_i \longrightarrow \alpha \bar{r}_i$$

↑
non-zero.

How?

$$\Psi(\{\bar{r}_i\}) \longrightarrow \Psi_\alpha(\{\bar{r}_i\}) = \alpha^{3N/2} \Psi(\{\alpha \bar{r}_i\})$$

Let's drop the "{ }":

$$\Psi(\bar{r}_i) \longrightarrow \Psi_\alpha(\bar{r}_i) = \alpha^{3N/2} \Psi(\alpha \bar{r}_i)$$

Preserve normalization on scaling

$$\therefore 1 = \int |\Psi_\alpha(\bar{r}_i)|^2 \prod_{i=1}^N d\bar{r}_i$$

Q: Does the expression for Ψ_α lead to a normalized Ψ_α ?

ie. $1 \stackrel{?}{=} \int |\alpha^{3N/2} \Psi(\alpha \bar{r}_i)|^2 \prod_{i=1}^N d\bar{r}_i$

$$\stackrel{?}{=} \alpha^{3N} \int |\Psi(\alpha \bar{r}_i)|^2 \prod_{i=1}^N d\bar{r}_i$$

$$\stackrel{?}{=} \int |\Psi(\alpha \bar{r}_i)|^2 \prod_{i=1}^N d(\alpha \bar{r}_i)$$

$d\bar{r}_i = dx_i dy_i dz_i$
3 of these

let $\bar{r}_i' = \alpha \bar{r}_i$

$$\stackrel{?}{=} \int |\Psi(\bar{r}_i')|^2 \prod_{i=1}^N d\bar{r}_i'$$

$= 1$ as Ψ is normalized. ✓

Uniform Scaling

$$\rho_{\alpha}(\bar{r}) = \alpha^3 \rho(\alpha \bar{r})$$

$$\rho_{\alpha}(\bar{r}) = N \int |\psi_{\alpha}(\bar{r}, \bar{r}_2, \dots, \bar{r}_N)|^2 \prod_{i=2}^N d\bar{r}_i$$

$$= N \int |\alpha^{3N/2} \psi(\alpha \bar{r}, \alpha \bar{r}_2, \dots, \alpha \bar{r}_N)|^2 \prod_{i=2}^N d\bar{r}_i$$

Remember to square it!

$$= \alpha^3 N \int |\psi(\alpha \bar{r}, \alpha \bar{r}_2, \dots, \alpha \bar{r}_N)|^2 \prod_{i=2}^N d(\alpha \bar{r}_i)$$

$$\text{Let } \bar{r}_2' = \alpha \bar{r}_2$$

$$\bar{r}_3' = \alpha \bar{r}_3 \text{ etc}$$

$$= \alpha^3 N \int |\psi(\alpha \bar{r}, \bar{r}_2', \bar{r}_3', \dots, \bar{r}_N')|^2 \prod_{i=2}^N d\bar{r}_i'$$

$$= \alpha^3 \rho(\alpha \bar{r})$$

Uniform Scaling

$$\langle \Psi_\alpha | T | \Psi_\alpha \rangle = \alpha^2 \langle \Psi | T | \Psi \rangle$$

$$\langle \Psi_\alpha | V_{ee} | \Psi_\alpha \rangle = \alpha \langle \Psi | V_{ee} | \Psi \rangle$$

$$\langle \Psi_\alpha | V_{en}(\bar{R}) | \Psi_\alpha \rangle = \alpha \langle \Psi | V_{en}(\alpha \bar{R}) | \Psi \rangle$$

$$\hat{T} = \sum_i -\frac{1}{2} \nabla_i^2$$

$$\hat{V}_{en} = \sum_{i=1}^N \frac{Z}{|\bar{r}_i - \bar{R}|}$$

$$\hat{V}_{ee} = \frac{1}{2} \sum_{ij} \frac{1}{r_{ij}}$$

can be generalized
to include more
nuclei.

* Try these out on one or 2 electron systems.

Consider the $\hat{V}_{ee}$ case:

$$\langle \Psi_\alpha | V_{ee} | \Psi_\alpha \rangle = \int \Psi_\alpha^*(\bar{r}_k) \left(\frac{1}{2} \sum_{ij} \frac{1}{r_{ij}} \right) \Psi_\alpha(\bar{r}_k) \prod_{k=1}^N d\bar{r}_k$$

$$= \int \alpha^{3N} \Psi^*(\alpha \bar{r}_k) \frac{1}{2} \sum_{ij} \frac{1}{r_{ij}} \Psi(\alpha \bar{r}_k) \prod d\bar{r}_k$$

$$= \int \Psi^*(\underline{\alpha \bar{r}_k}) \frac{1}{2} \sum_{ij} \frac{1}{\underline{r_{ij}}} \Psi(\underline{\alpha \bar{r}_k}) \prod d(\underline{\alpha \bar{r}_k})$$

needs to be $\underline{\alpha r_{ij}}$

$$r_{ij} = |\bar{r}_i - \bar{r}_j| = \frac{1}{\alpha} |\alpha \bar{r}_i - \alpha \bar{r}_j|$$

$$= \alpha \int \Psi^*(\alpha \bar{r}_k) \frac{1}{2} \sum_{ij} \frac{1}{|\alpha \bar{r}_i - \alpha \bar{r}_j|} \Psi(\alpha \bar{r}_k) \prod d(\alpha \bar{r}_k)$$

Now let $\bar{r}'_k = \alpha \bar{r}_k$

$$= \alpha \int \Psi^*(\bar{r}'_k) \frac{1}{2} \sum_{ii} \frac{1}{r_{ii}} \Psi(\bar{r}_k) \prod d\bar{r}'_k = \alpha \langle \Psi | V_{ee} | \Psi \rangle$$

Uniform Scaling

In DFT we will write quantities in terms of $\rho(\vec{r})$.

$$\therefore T[\rho] \equiv \langle \psi | T | \psi \rangle$$

~~$$J[\rho] + E_{xc}[\rho] \equiv \langle \psi | V_{ee} | \psi \rangle$$~~

The Coulomb part of $\langle \psi | V_{ee} | \psi \rangle$ will be $J[\rho]$ and the exchange part will be $E_x[\rho]$

We will use the scaling relations:

$$T[\rho_\alpha] = \alpha^2 T[\rho]$$

$$J[\rho_\alpha] = \alpha J[\rho]$$

$$E_x[\rho_\alpha] = \alpha E_x[\rho]$$